

## Surface modification of TiO<sub>2</sub> for visible light photocatalysis: Experimental and theoretical calculations of its electronic and optical properties

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Received 30 April 2019

Accepted 20 August 2019

Published 26 December 2019

Surface modification has been used as a method to create defects on TiO<sub>2</sub> materials, which can improve their desirable properties. In this paper, defected TiO<sub>2</sub> nano-powder was successfully synthesized by chemical reduction using NaBH<sub>4</sub> as the reducing agent at 300–400°C under argon atmosphere. High defect concentration can be produced by increasing process temperature. The modified TiO<sub>2</sub> shows good visible light absorption and photocatalytic activity on degradation of Rhodamine B (4–9 times higher than the pristine TiO<sub>2</sub>) with the visible light irradiation. Further XPS analysis and theoretical studies using full potential linearized augmented plane wave (FP-LAPW) method as implemented in wien2k code revealed the existence of oxygen vacancy and Ti<sup>3+</sup> in the modified samples. These types of defects were responsible for the modifications of the electronic and optical properties of TiO<sub>2</sub>, resulting in the improved photocatalytic activity in visible light irradiation.

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*Keywords:* NaBH<sub>4</sub> reduction; defect; density functional theory; TiO<sub>2</sub>; visible light absorption.

PACS number: 78.20.Bh

## 1. Introduction

TiO<sub>2</sub> is a wide bandgap semiconductor (3.0–3.2 eV) that can be used as photocatalyst.<sup>1</sup> It requires ultraviolet radiation (290–400 nm) to excite electrons from the valance band, not visible or fluorescent light,<sup>2</sup> which limit their application in utilizing major solar light spectrum. In order to enable light absorption in the visible region and to improve the photo-efficiency of TiO<sub>2</sub>, doping TiO<sub>2</sub> with nonmetals (such as C, N, S or F), metals (such as Fe, Cr or V) or metal oxides (such as V<sub>2</sub>O<sub>5</sub>, ZrO<sub>2</sub> or WO<sub>3</sub>) has been reported. Hybrid materials combining TiO<sub>2</sub> with carbon materials such as carbon nanotubes, activated carbons or TiO<sub>2</sub>-carbon microspheres were also investigated for the same propose.<sup>3–5</sup> In addition, new photocatalytic materials have been developed to achieve higher activities than TiO<sub>2</sub>-based photocatalysts; but the success cases for practical applications are few. One of the most viable and practical approaches in developing better photocatalysts is to use self-modify TiO<sub>2</sub> processes, including surface modification.

In this paper, surface modification has been conducted by chemical reduction using NaBH<sub>4</sub> to create defects in pristine TiO<sub>2</sub> powder. NaBH<sub>4</sub> is an inorganic compound that can be used as a versatile reducing agent. Thermal decomposition of NaBH<sub>4</sub> releases active hydrogen that is favorable for defect introduction to the surface of semiconductor.<sup>6</sup> The effect of temperature was studied in correlation to its optical and electronic properties. The full potential linearized augmented plane wave (FP-LAPW) method within density functional theory (DFT) was employed to investigate the effect of oxygen vacancy on the optical and electronic properties of the modified samples.

## 2. Experimental Details

### 2.1. Preparation of modified TiO<sub>2</sub> and sample characterization

Anatase TiO<sub>2</sub> (Sigma Aldrich Co. size <25 nm) and reduction agent NaBH<sub>4</sub> (ECP Ltd.) were used as the starting material. First, 0.75 g NaBH<sub>4</sub> was mixed with 2.0 g TiO<sub>2</sub>, heating then applied for 60 min in a furnace under Ar to 300–400°C (labelled X300, X350 and X400, cooled down, washed with DI water, and dried at 110°C.<sup>7</sup> The prepared samples were then characterized by XRD, UV-Vis-NIR and XPS.

### 2.2. Photocatalytic degradation of organic pollutant

Photocatalytic activity was evaluated by monitoring the decomposition of RhB in an aqueous solution under visible light irradiation (ULTRA VITALUX 300 W 230 V E27 with UV cut-off filter). Photocatalyst (50 mg) was mixed with RhB

solution (50 mL, 5 ppm), followed by 30 min stirring in dark to reach the adsorption equilibrium. Two milliliters were taken every time to filter by 0.2 mm PVDF syringe filter. The concentration of RhB was determined with a UV-Vis spectrophotometer by the peak intensity at 553 nm.

### 2.3. Method for theoretical calculations

The calculations were first carried out on the primitive unit cell of TiO<sub>2</sub> in anatase phase. The 2 × 2 × 2 96-atom anatase supercell was used to construct TiO<sub>2</sub> systems with oxygen vacancy  $V_o$  (concentration was set at 4.1%) and to calculate the  $V_o$  energy formation using FP-LAPW method as implemented in wien2k code.<sup>8</sup> The wavefunction for core states is taken as atomic like function, while for the valence states, it is considered as plane waves. The muffin-tin radius for core states is denoted as  $R_{MT}$ . Outside  $R_{MT}$ , valence states are accounted. In this work, nonoverlapping  $R_{MT}$  for Ti and O atoms were 1.93 and 1.75 a.u, respectively. Energy convergence criteria were set at 10<sup>-4</sup> eV. For calculation of electronic and optical properties, a  $k$  mesh of 36 points was sampled in the irreducible Brillouin zone (IBZ). The maximum cut-off for the plane wave  $k_{max}$  is 6.5. Exchange correlations effects are considered under generalized gradient approximation (GGA), GGA+U (with U as Hubbard potential, set as 10.303 eV) and Tran–Blaha modified Becke–Johnson (TB-mBJ) approximation.

## 3. Results and Discussion

### 3.1. Structures features and physical properties

As shown in Fig. 1(a), small changes such as peak boardening and reduction of the peak intensity were observed in the XRD peaks of the modified TiO<sub>2</sub> in comparison to the pristine one. The diffraction peaks of all samples were assigned for TiO<sub>2</sub> anatase (JCPDS card 78-2486) indicating that all samples have the same crystalline structure (anatase). Only peak intensity was reduced gradually with increasing. The deformation of cystalline lattice may occur and defect be created, thus resulting in low peak intensity.<sup>7</sup> Table 1 shows the crystallite size and lattice strain of the pristine and NaBH<sub>4</sub> modified TiO<sub>2</sub> estimated by Debye–Scherrer and Stokes–Wilson equations. As the temperature of reduction process gets higher, the crystallite size slightly reduces (from 63.5 to 48.8 nm), and the lattice strain slightly increases (0.010–0.012). Strain broadening is often triggered by lattice imperfection, including vacancies, interstitials and substitutions.<sup>9</sup>

Figure 1(b) depicts the optical properties of pristine and modified TiO<sub>2</sub> samples which are represented by the light absorption behavior. The absorbance intensity from 400 to 800 nm increased as the NaBH<sub>4</sub> treatment temperature increased. All modified samples were able to absorb visible light while pristine TiO<sub>2</sub> has no absorption. The enhancement of light visible light absorption in TiO<sub>2</sub> were also reported in Refs. 6 and 10–13.

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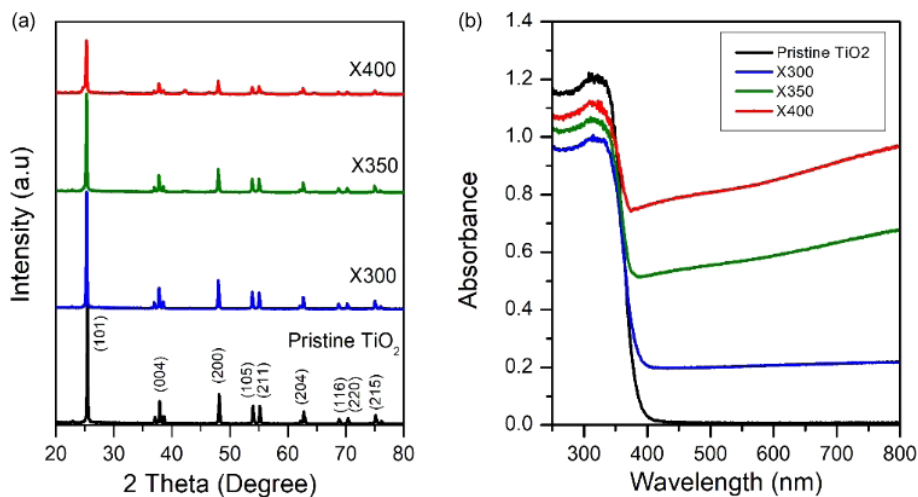


Fig. 1. (Color online) (a) XRD pattern and (b) UV-Vis spectrum of pristine and  $\text{NaBH}_4$  modified  $\text{TiO}_2$ .

Table 1. Crystallite size and lattice strain of pristine and  $\text{NaBH}_4$  modified  $\text{TiO}_2$ .

Samples	FWHM	Peak position	Crystallite size (nm)	Lattice strain
$\text{TiO}_2$ pristine	0.128	25.41	63.5	0.00991
X300	0.126	25.22	61.6	0.00981
X350	0.134	25.26	60.2	0.01044
X400	0.157	25.24	48.8	0.01224

Table 2. Percentage area of O 1s.

Sample	O 1s (Ti-O) (%)	O 1s (O-H) (%)
Pristine $\text{TiO}_2$	68.10	31.90
X300	62.65	37.35
X350	56.26	43.74
X400	44.54	55.46

Figure 2 compares the XPS spectra of Ti 2p of pristine and modified  $\text{TiO}_2$  at different temperatures.  $\text{Ti}^{4+}$  and  $\text{Ti}^{2+}$  were detected in all samples, representing the same state of Ti atoms that existed in these samples. However, additional peak of  $\text{Ti}^{3+}$  was detected in X350 [Fig. 2(c)], indicating the existence of defects. Hydroxide (O-H) and Ti (Ti-O) in the oxygen state (O 1s) were also detected. The presence of hydroxide is predictable as the defect formation involves the attraction of active hydrogen with the oxygen in the lattice structures, thus it can be associated with the surface defects and oxygen vacancies.<sup>14</sup> Table 2 shows the percentage of O 1s (H-O) in the sample, which increases gradually with increasing temperature. High percentage of O 1s (H-O) indicates the existence of distortion on the  $\text{TiO}_2$  surface and more severe conditions compared to the distortion in bulk.<sup>7</sup>

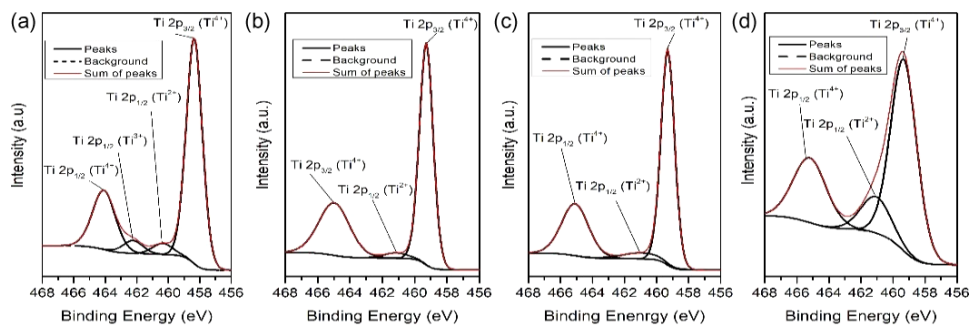


Fig. 2. (Color online) XPS spectra of Ti 2p: (a) pristine  $\text{TiO}_2$ , (b) X300, (c) X350 and (d) X400.

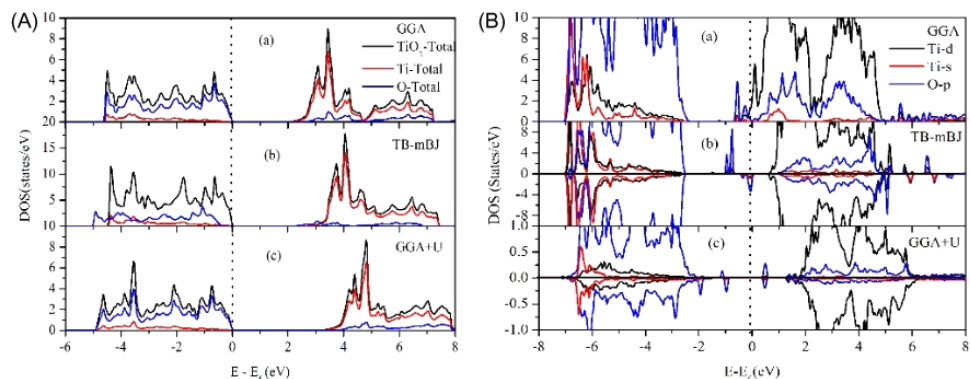


Fig. 3. (Color online) Density of states of (a) pristine and (b) defected  $\text{TiO}_2$  with oxygen vacancy  $V_o$ .

XPS also indicate the electronic band position of the materials (valence band). The valence band maximum (VBM) is estimated by a linear extrapolation of the peak to the baseline.<sup>15</sup> The pristine  $\text{TiO}_2$  have a band edge position of 1.7 eV below the Fermi energy. Meanwhile, the band edge of modified  $\text{TiO}_2$  were slightly lower by 0.1; 0.4 and 0.2 eV for X300; X350 and X400, respectively. There is a possibility of the formation of an extra band between the Fermi level (0 eV) and the top of the VB triggered by the defect of  $\text{Ti}^{3+}$  and oxygen vacancies.<sup>15</sup> To confirm the formation of extra band of the modified  $\text{TiO}_2$ , the theoretical calculation based on FP-LAPW DFT was conducted and discussed in Sec. 3.3.

### 3.2. Photocatalytic activity

Table 3 presents the kinetic constant of pristine and modified  $\text{TiO}_2$  on RhB degradation under visible light. Sample X350 has the highest degradation rate of  $k = 0.0018/\text{min}$ , followed by X300 and X400 with  $k = 0.0011$  and  $0.0008/\text{min}$ , respectively. The pristine  $\text{TiO}_2$  has the lowest degradation rate as it does not absorb visible light. As X400 has higher oxygen vacancies O 1s (H-O), it has lower

Table 3. Kinetics constant of RhB degradation under visible light.

Catalyst	Kinetics constant ( $\text{min}^{-1}$ )	$R^2$
Pristine $\text{TiO}_2$	0.0002	0.9601
X300	0.0011	0.9648
X350	0.0018	0.9811
X400	0.0008	0.9864

photocatalytic activity compared to X300 and X350. This can be explained by the possibility of samples having different type and concentration of defects. By decreasing the relative defect concentration of bulk to surface in the  $\text{TiO}_2$ , the photocatalytic efficiency is significantly enhanced.<sup>14,16</sup>

### 3.3. Theoretical studies

The calculated bandgap energy for pristine  $\text{TiO}_2$  by using GGA, GGA+U and TB-mBJ approximations is 2.015, 3.246 and 2.872 eV, respectively. Figure 3(a) shows the total density of states for anatase  $\text{TiO}_2$  unit cell. Based on Fig. 3(a), the conduction band is made up of Ti-3d valence states while O-2p states contribute in the formation of the valence band. The supercell model, oxygen vacancy  $V_o$ , was created; and the effect of  $V_o$  on the bandgap and state of density of states was investigated. Density of states (DOS) of Ti-d, Ti-s and O-p states are plotted [Fig. 3(b)]. It shows the new localized states between the conduction band and valence band are O-2p states which are created due to the  $V_o$  defect. TB-mBJ gives bandgap  $\sim 0.7$  eV in spin up channel and  $\sim 0.446$  eV in spin down channels. The conduction band and valence band both shift to the lower energy levels and become narrower. DOS shows that new O-2p states are located at 1.1 eV above the VBM due to  $V_o$  defect. Similar to GGA, these localized energy states in the valence band act as the trap centers for photoexcited holes.

GGA+U approximation depicts a different picture as compared to GGA and TB-mBJ schemes. The bandgap is improved and gives  $\sim 1.5$  eV in spin down channel whereas it reduced to  $\sim 0.426$  eV in the spin up channel. The energy states due to  $V_o$  are located in the conduction band and valence band. These are O-2p states which form shallow levels at the CBM and localized energy states at 1.5 eV below the conduction band. In the valence band, they form shallow donor energy levels at VBM edge and localized energy states at  $\sim 0.7$  eV above the VBM. The possible explanation of this behavior is that  $V_o$  acts like donor impurity with a charge state +2. With the creation of  $V_o$ , three O-2p states but four electrons were removed, which are responsible for the valence band formation. Now, two Ti-3d electrons need to transfer from the conduction band to valence band to fill the O-2p<sub>y</sub> and O-2p<sub>z</sub> states in the pristine  $\text{TiO}_2$  unit cell, residing in the conduction band.

These unpaired Ti-3d electrons behave like free electrons in the conduction band, and are responsible for the *n*-type conductivity in  $\text{TiO}_2$ . Reports cited the effect

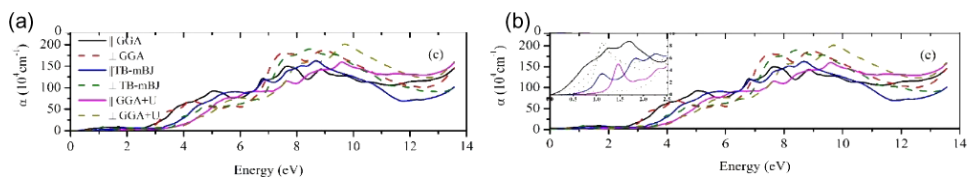


Fig. 4. (Color online) Absorption coefficient spectra of (a) pristine and (b) defected TiO<sub>2</sub>.

of  $V_o$  defects on the photocatalytic activity.<sup>17</sup> The presence of oxygen vacancies in TiO<sub>2</sub> can increase or reduce the photocatalytic activity depending on the factors of the concentration of  $V_o$ , oxygen vacancy at the surface or in the bulk. They provide trap centers for electrons and holes, either increase or decrease the electron–hole recombination rate. It is also reported that the electrons can easily be photo-excited to these localized states by visible light absorption and cause an increase in the photocatalytic activity.<sup>17</sup>

The threshold energy for absorption is closely related to the optical bandgap of the material. For pristine TiO<sub>2</sub>, it is calculated as 2.3–2.5, 3.2–3.5 and 3.5–4.0 eV in parallel and perpendicular directions of incident photons with GGA, TB-mBJ and GGA+U approximations, respectively [Fig. 4(a)]. Figure 4(b) shows the light absorption in the visible energy range of  $\sim 0.5$ –2.5 eV for defected TiO<sub>2</sub> and confirmed by the experimental results presented in Fig. 2. Although the absorption in visible light is very low compared to the UV region, the photocatalytic activity is significantly enhanced by  $V_o$  defects, as shown in Table 3.

#### 4. Conclusion

Defects have been successfully introduced to the pristine TiO<sub>2</sub> via NaBH<sub>4</sub> treatment. Its concentration can be controlled by temperature. The modified TiO<sub>2</sub> has good visible light absorption and high photocatalytic activity towards RhB degradation (4–9 times higher than pristine TiO<sub>2</sub>). This may come from the formation of  $V_o$  and Ti<sup>3+</sup>. Theoretical calculations confirm the presence of extra energy states created by  $V_o$ . The bandgap is reduced in the TiO<sub>2</sub> with  $V_o$  defects as compared to pristine TiO<sub>2</sub>. These changes lead to a disordered layer and a broad visible light absorption which is in line with the experimental findings.

#### References

1. A. Fujishima and K. Honda, *Nature* **238**, 37 (1972).
2. X. Wang *et al.*, *J. Hazard. Mater.* **262**, 16 (2013).
3. M. Sillanpää, M. C. Ncibi and A. Matilainen, *J. Environ. Manage* **208**, 56 (2018).
4. A. Fujishima, X. Zhang and D. A. Tryk, *Surf. Sci. Rep.* **63**, 515 (2008).
5. X. Chen, L. Liu and F. Huang. *Chem. Soc. Rev.* **44**, 1861 (2015).
6. H. Tan *et al.*, *Nanoscale* **6**, 10216 (2014).
7. D. Ariyanti *et al.*, *Mater. Chem. Phys.* **199**, 571 (2017).

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8. P. K. S. Blaha *et al.*, WIEN2k: An augmented plane wave plus local orbitals program for calculating crystal properties (2001).
9. D. Balzar, *J. Res. Natl. Inst. Stand. Technol.* **98**, 321 (1993).
10. J. Dong *et al.*, *ACS Appl. Mater. Interfaces* **6**, 1385 (2014).
11. H. Wang *et al.*, *Catal. Commun.* **60**, 55 (2015).
12. G. Zhu *et al.*, *ChemCatChem.* **7**, 2614 (2015).
13. T. Lin *et al.*, *Energy Environ. Sci.* **7**, 967 (2014).
14. J. Yan *et al.*, *Phys. Chem. Chem. Phys.* **15**, 10978 (2013).
15. S. Pan *et al.*, *J. Mater. Chem. A* **3**, 11437 (2015).
16. M. Kong *et al.*, *J. Am. Chem. Soc.* **133**, 16414 (2011).
17. X. Pan *et al.*, *Nanoscale* **5**, 3601 (2013).